

# 2025.01.27

鈴木 貴(大阪大学)

## 1. 量子化する爆発機構

Model

Non-equilibrium thermo-dynamics

# $\Omega \subset \mathbf{R}^2$ bounded domain, $\partial \Omega$ smooth

1. Smoluchowski Part

$$\begin{aligned} u_t &= \nabla \cdot (\nabla u - u \nabla v) \\ \frac{\partial u}{\partial \nu} - u \frac{\partial v}{\partial \nu} \Big|_{\partial \Omega} &= 0, \ u|_{t=0} = u_0(x) > 0 \end{aligned}$$

1. total mass conservation  $\frac{d}{dt} ||u(t)||_1 = 0$ 

2. free energy decreasing

2. Poisson Part

Green's function

A = U - TS

 $-\Delta v = u, \ v|_{\partial \Omega} = 0$ 

G(x, x') = G(x', x)

 $egin{aligned} u &= u(x,t) \geq 0 & ext{density} \ j &= abla u + u 
abla v & ext{flux} \ _{ ext{diffusion mass velocity}} & u_t + 
abla \cdot j &= 0 & ext{conservation law} \ v &= (-\Delta)^{-1} u & ext{potential} \end{aligned}$ 

attractive (chemotaxis, gravitation) action at a distance (long range potential) symmetry (action-reaction)

 $\begin{aligned} & \text{scaling} \\ & u_{\mu}(x,t) = \mu^2 u(\mu x, \mu^2 t), \ \mu > 0 \\ \hline & u_{\mu}(x) = \mu^2 u(\mu x), \ \mu > 0 \\ \hline & \|u\|_1 = \|u_{\mu}\|_1 \equiv \lambda \end{aligned} \Leftrightarrow \ n = 2 \quad \text{critical dimension} \\ F(u) = \int_{\mathbf{R}^2} u(\log u - 1) - \frac{1}{2} \langle \Gamma * u, u \rangle, \ \Gamma(x) = \frac{1}{2\pi} \log \frac{1}{|x|} \\ F(u_{\mu}) = \left(2\lambda - \frac{\lambda^2}{4\pi}\right) \log \mu + \mathcal{F}(u) \quad \text{critical mass} \quad \lambda = 8\pi \end{aligned}$ 

Boltzmann equation

Boltzmann Poisson equation ~ stationary state

$$\nabla = \left(\begin{array}{c} \frac{\partial}{\partial x_1} \\ \frac{\partial}{\partial x_2} \end{array}\right)$$

Euler's equation of motion

 $v_t + (v \cdot \nabla)v = -\nabla p, \ \nabla \cdot v = 0, \ \nu \cdot v|_{\partial\Omega} = 0$ 

2D 
$$\omega = \nabla^{\perp} \cdot v \rightarrow \omega_t + \nabla \cdot (v\omega) = 0, \ \nabla \cdot v = 0$$
  
 $v = \nabla^{\perp} \psi$  stream function  $\nabla^{\perp} = \begin{pmatrix} -\frac{\partial}{\partial x_2} \\ \frac{\partial}{\partial x_1} \end{pmatrix}$ 

$$\omega_t + \nabla \cdot (\omega \nabla^\perp \psi) = 0, \ -\Delta \psi = \omega$$
 vorticity equation  
boundary condition  $\longrightarrow \psi|_{\partial\Omega} = 0$ 

Green function

$$\begin{split} -\Delta_x G(x, x') &= \delta_{x'}(dx), \ G(x, x')|_{x \in \partial \Omega} = 0 \quad \rightarrow \\ \omega_t + \nabla \cdot (\omega \nabla^\perp \psi) &= 0, \ \psi(\cdot, t) = \int_{\Omega} G(\cdot, x') \omega(x', t) dx' \\ G(x, x') &= G(x', x) \text{ action reaction law} \end{split}$$

$$\frac{dx_i}{dt} = \nabla_{x_i}^{\perp} H_\ell$$

point vortex Hamiltonian  

$$H_{\ell}(x_1, \dots, x_{\ell}) = \sum_{i} \frac{\alpha_i^2}{2} R(x_j) + \sum_{i < j} \alpha_i \alpha_j G(x_i, x_j)$$
Robin function 
$$R(x) = \left[ G(x, x') + \frac{1}{2\pi} \log |x - x'| \right]_{x' = x}$$

Thermal equilibrium

Gibbs theory of statistical mechanics

#### micro-canonical statistics

$$\frac{dq_i}{dt} = \frac{\partial H}{\partial p_i}, \ \frac{dp_i}{dt} = -\frac{\partial H}{\partial q_i}, \ 1 \le i \le \ell$$

$$\mathbf{R}^{6\ell}/\{H = E\}$$
  

$$x = (q_1, \dots, q_\ell, p_1, \dots, p_\ell), \ dx = dE \cdot \frac{d\Sigma(E)}{|\nabla H|}$$
  

$$d\Sigma(E) \ \leftrightarrow \ \left\{x \in \mathbf{R}^{6\ell} \mid H(x) = E\right\}$$

micro-canonical measure

weight factor

$$d\mu^{E,N} = \frac{1}{W(E)} \cdot \frac{d\Sigma(E)}{|\nabla H|}, \ W(E) = \int_{\{H=E\}} \frac{d\Sigma(E)}{|\nabla H|}$$

## canonical statistics

Boltzmann constant 
$$\mathbf{R}^{6N}/\{T\}, \ \beta = 1/(kT) \quad \text{inverse temperature}$$
$$d\mu^{\beta,N} = \frac{e^{-\beta H}dx}{Z(\beta,N)}, \ Z(\beta,N) = \int_{\mathbf{R}^{6N}} e^{-\beta H}dx$$
canonical measure

Boltzmann-Poisson equation

thermo-dynamical relation

$$\beta = \frac{\partial}{\partial E} \log W(E)$$

order structure in negative temperature

$$H_{\ell}(x_{1},\ldots,x_{\ell}) = \sum_{i} \frac{\alpha_{i}^{2}}{2} R(x_{j}) + \sum_{i < j} \alpha_{i} \alpha_{j} G(x_{i},x_{j}) \quad \text{principle of equal probability} \quad \longleftrightarrow \quad -\Delta v = u, \quad v|_{\partial\Omega} = 0$$
  
$$\alpha_{i} = \hat{\alpha}, \ \hat{\alpha}\ell = 1, \ \hat{H}_{\ell} = H, \ \hat{\alpha}^{2}\ell\hat{\beta} = \beta \qquad \ell \uparrow +\infty \qquad \qquad u = \frac{\lambda e^{v}}{\int_{\Omega} e^{v} dx}, \ \lambda = \|u\|_{1}$$

## Hamiltonian is recursive

$$u = \frac{\lambda e^v}{\int_{\Omega} e^v dx} \qquad -\Delta v = u, \ v|_{\partial\Omega} = 0$$

Boltzmann

$$\rightarrow \quad -\Delta v = \frac{\lambda e^v}{\int_\Omega e^v}, \ v|_{\partial\Omega} = 0$$

Nonlinear eigenvalue problem





point vortices ~ negative inverse temperature L. Onsager 49

$$G(x, x') = G(x', x)$$
  

$$R(x) = \left[G(x, x') + \frac{1}{2\pi} \log |x - x'|\right]_{x' = x}$$

Theorem A (Nagasaki-S. 90)  

$$\{(\lambda_k, v_k)\}, \ \lambda_k \to \lambda_0 \in (0, \infty), \ \|v_k\|_{\infty} \to \infty$$
  
 $\rightarrow \lambda_0 = 8\pi\ell, \ \ell \in \mathbf{N}, \ \exists S \subset \Omega, \ \sharp S = \ell$  stationary quantization

(sub-sequence)  $v_k o v_0$  loc. unif. in  $\overline{\Omega} \setminus \mathcal{S}$ 

$$v_0(x) = 8\pi \sum_{x_0 \in \mathcal{S}} G(x, x_0), \ \mathcal{S} = \{x_1^*, \dots, x_\ell^*\}$$

$$\nabla H_{\ell}|_{(x_1,\dots,x_{\ell})=(x_1^*,\dots,x_{\ell}^*)} = 0$$
  
$$H_{\ell}(x_1,\dots,x_{\ell}) = \frac{1}{2} \sum_i R(x_i) + \sum_{i < j} G(x_i,x_j)$$



$$\begin{split} \hline \text{Geometric background} \\ -\Delta v &= \sigma e^v, \ \sigma = \frac{\lambda}{\int_{\Omega} e^v dx} \\ \Leftrightarrow \ \exists F = F(z), \ z \in \Omega \subset \mathbf{R}^2 \cong \mathbf{C} \quad \text{meromorphic function} \\ \rho(F) &= \left(\frac{\sigma}{8}\right)^{1/2} e^{v/2} = \frac{|F'|}{1+|F|^2} \quad \text{spherical derivative} \\ -\Delta v &= \sigma e^v, \ v|_{\partial\Omega} = 0 \ \Leftrightarrow \ \rho(F)|_{\partial\Omega} = \left(\frac{\sigma}{8}\right)^{1/2} \end{split}$$

$$\begin{split} \hat{F} &= \sqrt{8} \circ F : \Omega \to S^2 \quad \text{ conformal} \\ \left. \frac{d\Sigma}{ds} \right|_{\partial \Omega} &= \sigma^{1/2} \quad & (S^2, d\Sigma) \text{ round sphere} \\ \left. |S^2| &= 8\pi \end{split}$$

$$\int_{\partial\Omega} \frac{d\Sigma}{ds} ds = \left|\partial\Omega\right| \sigma^{1/2}$$

immersed length of  $\hat{F}(\partial \Omega)$ 

 $\hat{F}$ 

# **Proof of Theorem A**

- 1. Liouville integral
- 2. boundary reflection
- 3. elliptic regularity
- 4. complex function theory
- 4-1. maximum principle
- 4-2. Montel's theorem
- 4-3. theorem of coincidence
- 4-4. residue analysis

$$\int_{\Omega} \left(\frac{d\Sigma}{ds}\right)^2 dx = 8 \int_{\Omega} \rho(F)^2 dx = \int_{\Omega} \sigma e^{v}$$
  
immersed area of  $\hat{F}(\Omega)$ 

$$\lambda = \int_{\Omega} \sigma e^v \to 8\pi\ell$$

 $\Leftrightarrow$  total mass quantization due to  $\ell$ -covering

#### Blowup analysis

$$\begin{split} &\Omega \subset \mathbf{R}^2 \text{: open set, } V \in C(\overline{\Omega}) \\ &-\Delta v = V(x)e^v, \ 0 \leq V(x) \leq b \quad \text{in } \Omega \\ &\int_{\Omega} e^v \leq C \end{split}$$

## Theorem a (Li-Shafrir 94)

 $\{(V_k, v_k)\}$  solution sequence  $V_k \to V$  loc. unif. in  $\Omega$ 

- $\Rightarrow$   $^\exists$  sub-sequence with the alternatives;
- 1.  $\{v_k\}$ : loc. unif. bdd in  $\Omega$

2. 
$$\exists S \subset \Omega, \ \ \ S < +\infty$$
  
 $v_k \to -\infty$  loc. unif. in  $\Omega \setminus S$   
 $S = \{x_0 \in \Omega \mid \exists x_k \to x_0, \ v_k(x_k) \to +\infty$   
 $V_k(x)e^{v_k}dx \rightharpoonup \sum_{x_0 \in S} m(x_0)\delta_{x_0}(dx)$  in  $\mathcal{M}(\Omega)$   
 $m(x_0) \in 8\pi \mathbf{N}$ 

3.  $v_k \to -\infty$  loc. unif. in  $\Omega$ 

#### Comments

1. mass quantization for variable coefficients without boundary condition

- 2. possible collapse collision
- 3. many applications together with the proof

<u>prescaled analysis</u> ...Brezis-Merle 91 linear theory ⇒

1, 2 with  $m(x_0) \ge 4\pi$  (rough estimate), 3

```
2... localized to B = B(0, R)

-\Delta v_k = V_k(x)e^{v_k}, V_k(x) \ge 0 \text{ in } B

V_k \to V \text{ unif. in } \overline{B}, \max_{\overline{B}} v_k \to +\infty

\max_{\overline{B} \setminus B_r} v_k \to -\infty, \forall r \in (0, R)

\lim_k \int_B V_k e^{v_k} = \alpha, \int_B e^{v_k} \le C

\Rightarrow \alpha \in 8\pi \mathbf{N}
```

#### Scaling

 $v_k(x_k) = \|v_k\|_{\infty}, \ x_k \to 0$   $\tilde{v}_k(x) = v_k(\delta_k x + x_k) + 2\log \delta_k, \ \delta_k = e^{-v_k(x_k)/2} \to 0$   $-\Delta \tilde{v}_k = V_k(\delta_k x + x_k)e^{\tilde{v}_k}, \ \tilde{v}_k \le \tilde{v}_k(0) = 0 \text{ in } B(0, R/2\delta_k)$   $\int_{B(0, R/2\delta_k)} e^{\tilde{v}_k} \le C$ 

#### Liouville property

Theorem b (Chen-Li 1991) Liouville property  $-\Delta v = e^{v} \text{ in } \mathbf{R}^{2}, \ \int_{\mathbf{R}^{2}} e^{v} < +\infty$   $\overrightarrow{}$   $v(x) = \log \left\{ \frac{8\mu^{2}}{(1+\mu^{2}|x-x_{0}|^{2})^{2}} \right\}, \ x_{0} \in \mathbf{R}^{2}, \ \mu > 0$   $\int_{\mathbf{R}^{2}} e^{v} = 8\pi$  Tail control in the scaled variable

#### Theorem c (Shafrir 91)

 $-\Delta v = V(x)e^v, \ 0 < a \le V(x) \le b \text{ in } \Omega$ 

 $K \subset \Omega \quad \underset{K}{\operatorname{compact}} \longrightarrow$  $\sup_{K} v + c_1 \inf_{\Omega} v \leq \exists c_2, \ \exists c_1 \geq 1$ 

residual vanishing

Simplicity

Theorem d (Y.Y. Li 99)  $\max_{\partial B} v_k - \min_{\partial B} v_k \le C, \ \|\nabla V_k\|_{\infty} \le C$   $\longrightarrow \alpha = 8\pi$   $\left| v_k(x) - \log \frac{e^{v_k(0)}}{\left(1 + \frac{V_k(0)}{8}e^{v_k(0)}|x|^2\right)^2} \right| \le C, \ x \in B$  1. non-radial bifurcation on annulus (S.S. Lin 89 Nagasaki-S. 90b)

- 2. effective bound of blowup points for simply-connected domain (S.-Nagasaki 89 Grossi-F.Takahashi 10)
- 3. classification of singular limits (Nagasaki-S. 90a)
- 4. spherical mean value theorem (S. 90)
- 5. localization (Brezis-Merle 91)
- 6. entire solution (W. Chen-C. Li 91)
- 7.  $\sup + \inf \operatorname{inequality} (\operatorname{Shafrir} 92)$

8. uniqueness (S. 92)

9. field-particle duality (S. 92 Wolansky 92)

10. singular perturbation (Weston 78Moseley 83 S. 93 Baraket-Pacard 98Esposito-Grossi-Pistoia 05

del Pino-Kowarzyk-Musso 05)

- 11. blowup analysis (Li-Shafrir 94)
- 12. Chern-Simons theory (Tarantello 96)
- 13. global bifurcation (S.-Nagasaki 89 Mizoguchi-S. 97 Chang-Chen-Lin 03)
- 14. min-max solution (Ding-Jost-Li Wang 99)
- 15. local uniform esitmate (Y.Y. Li 99)
- 16. variable coefficient (Ma-Wei 01)
- 17. refined asymptotics (Chen-Lin 02)
- 18. topological degree (Li 99 C.C. Chen-C.S. Lin 03 Malchiodi 08)
- 19. asymptotic non-degeneracy (Gladiali-Grossi 04 Grossi-Ohtsuka-S. 11)
- 20. isoperimetric profile (Lin-Lucia 06)
- 21. deformation lemma (Lucia 07)
- 22. Morse index (Gladiali-Grossi 09)



$$-\Delta v = \frac{\lambda e^v}{\int_{\Omega} e^v} \text{ in } \Omega \subset \mathbf{R}^2$$
$$v|_{\partial \Omega} = 0$$

#### 2. 動的な階層の循環

From quasi-equilibrium to equilibrium Staniscia-Chavanis-Ninno-Fanelli 09 0.45 state of the system equilibrium 0.4 0.35 Poisson 0.3  $-\Delta v = u, \ v|_{\partial\Omega}$ 0.25 = 0Σ 0.2 0.15 kinetic 0.1 0.05 time 0 106 100 10<sup>2</sup>  $10^{3}$ 10<sup>4</sup> 10<sup>5</sup>  $10^{1}$ initial quasi-equilibrium relaxation relaxation static

Chavanis 08 relaxation to the equilibrium in the point vortices, kinetic equation + maximum entropy production Sire-Chavanis 02 motion of the mean field of many self-gravitating Brownian particles, BBGKY hierarchy + factorization

$$\left| \begin{array}{c} u_t = \Delta u - \nabla \cdot u \nabla v, \quad \frac{\partial u}{\partial \nu} - u \frac{\partial v}{\partial \nu} \right|_{\partial \Omega} = 0 \end{array} \right|_{\text{Boltzmann}}^{\text{Smoluchowski}} \left| \begin{array}{c} u = \frac{\lambda e^v}{\int_{\Omega} e^v dx}, \quad \lambda = \|u\|_1 \\ \end{array} \right|_{\Omega} \right|_{\Omega} = 0 \left| \begin{array}{c} \sum_{x \in V} e^{x} dx \\ \sum_{x \in V} e^{x} dx \\ \sum_{x \in V} e^{x} dx \\ \end{array} \right|_{\Omega} \left| \begin{array}{c} \sum_{x \in V} e^{x} dx \\ \sum_{x \in V} e^{x} dx \\ \sum_{x \in V} e^{x} dx \\ \end{array} \right|_{\Omega} \left| \begin{array}{c} \sum_{x \in V} e^{x} dx \\ \sum_{x \in V} e^{$$

Hamilton system of many particles with inner interaction of long range

## Results

$$\begin{split} \Omega &\subset \mathbf{R}^2 \text{ bounded domain, } \partial\Omega \text{ smooth} \\ \text{1. Smoluchowski Part} \\ u_t &= \nabla \cdot (\nabla u - u \nabla v) \\ \left. \frac{\partial u}{\partial \nu} - u \frac{\partial v}{\partial \nu} \right|_{\partial \Omega} = 0, \ \left. u \right|_{t=0} = u_0(x) > 0 \end{split}$$

2. Poisson Part

 $-\Delta v = u, \quad v|_{\partial\Omega} = 0$ 

Theorem B 
$$T < +\infty \longrightarrow$$
  
 $u(x,t)dx \rightharpoonup \sum_{x_0 \in S} m(x_0)\delta_{x_0}(dx) + f(x)dx$ 

 $m(x_0) \in 8\pi {f N}$  collapse mass quantization possibly with sub-collapse collision

#### blowup set

$$\begin{split} \mathcal{S} &= \{ x_0 \in \overline{\Omega} \mid \exists x_k \to x_0, \ t_k \uparrow T, \ u(x_k, t_k) \to +\infty \} \subset \Omega & \ \begin{tabular}{ll} \mathsf{Corol} \\ & \sharp \mathcal{S} < +\infty & \mbox{finiteness of blowup points} \\ & 0 < f = f(x) \in L^1(\Omega) \cap C(\overline{\Omega} \setminus \mathcal{S}) & \mbox{measure theoretic regular part} \\ & \ \begin{tabular}{ll} \mathsf{c.f. Gross} \\ & \ \mathsf{c.f. Gross} \\ & \ \mathsf{c.f. Gross} \\ \end{tabular} \end{split}$$

#### Theorem C

$$T = +\infty, \lim_{t \uparrow +\infty} \sup \|u(\cdot, t)\|_{\infty} = +\infty$$
  
$$\longrightarrow \quad \lambda \equiv \|u_0\|_1 = 8\pi\ell, \ \exists \ell \in \mathbf{N} \quad \text{initial mass quantization}$$
  
$$\exists x_* \in \Omega^\ell \setminus D, \ \nabla H_\ell(x_*) = 0 \quad \text{recursive hierarchy}$$

point vortex Hamiltonian  

$$H_{\ell}(x_1, \cdots, x_{\ell}) = \frac{1}{2} \sum_{j}^{\text{Robin function}} R(x_j) + \sum_{i < j}^{\text{Green function}} G(x_i, x_j)$$

Corollary 1  $T < +\infty$  if  $\nexists$  stationary solution or  $\mathcal{F}(u_0) \ll -1$ 

and 
$$\lambda \notin 8\pi \mathbf{N}$$
 or  
 $\lambda \in 8\pi \ell, \ \ell \in \mathbf{N}, \ \not\exists \text{ critical point of } H_{\ell}$ 

$$\begin{array}{c} \begin{array}{c} \text{exclusion of boundary blowup} \\ t_k) \to +\infty \} \subset \Omega \end{array} \quad \begin{array}{c} \begin{array}{c} \text{Corollary 2} \\ \end{array} & \Omega \\ \end{array} \quad \begin{array}{c} \Omega \\ \end{array} \quad \begin{array}{c} \text{convex} \\ \lambda \neq 8\pi \\ \end{array} \\ \end{array} \\ \begin{array}{c} \Rightarrow \\ T < +\infty \\ \end{array} \quad \begin{array}{c} \text{or} \\ T = +\infty \\ \end{array} \\ \begin{array}{c} \text{pre-compact orbit} \\ \end{array} \\ \end{array} \\ \end{array} \\ \begin{array}{c} \exists \\ \text{stationary solution} \end{array} \end{array}$$

Preliminary  $\sim$  Green's function and the weak form

action reaction law

$$\begin{aligned} G(x,x') &= \Gamma(x-x') + K(x,x'), \ K = K(x,x') \in C^{1+\theta,\theta}(\Omega \times \overline{\Omega}) \cap C^{\theta,1+\theta}(\overline{\Omega} \times \Omega) \\ x_0 &\in \partial\Omega \longrightarrow G(x,x') = E(X,X') + K(x,x'), \ K = K(x,x') \in (C^{1+\theta,\theta} \cap C^{\theta,1+\theta})(\overline{\Omega \cap B(x_0,R)} \times \overline{\Omega \cap B(x_0,R)}) \end{aligned}$$

$$X: \overline{\Omega \cap B(x_0, 2R)} \to \overline{\mathbf{R}_+}^2 = \{ (X_1, X_2) \mid X_2 \ge 0 \} \qquad X(\partial \Omega \cap B(x_0, 2R)) \subset \partial \mathbf{R}_+^2$$
$$E(X, X') = \Gamma(X - X') - \Gamma(X - X'_*) \qquad X = \begin{pmatrix} X_1 \\ X_2 \end{pmatrix} \mapsto X_* = \begin{pmatrix} X_1 \\ -X_2 \end{pmatrix}$$

# Proof of Theorem B

$$\begin{array}{|c|c|} \varepsilon \text{ regularity} \end{array} \begin{array}{|c|c|} \frac{1}{p+1} \frac{d}{dt} \|u\|_{p+1}^{p+1} = -\frac{4p}{(p+1)^2} \|\nabla u^{\frac{p+1}{2}}\|_2^2 + \frac{p}{p+1} \|u\|_{p+3}^{p+3} & \|z\|_{p+1}^{p+1} \le C_p \|z\|_1 \|z\|_{H^1}^p \\ n = 2 & \text{Gagliardo-Nirenberg inequality} \\ \exists \varepsilon_0 > 0, \ \|u_0\|_1 < \varepsilon_0 \ \Rightarrow \ T = +\infty, \ \|u(\cdot, t)\|_{\infty} \le C & \text{Gagliardo-Nirenberg inequality} \\ \end{array}$$

Moser's iteration scheme maximal regularity

## localization

$$\lim_{R \downarrow 0} \limsup_{t \uparrow T} \|u(\cdot, t)\|_{L^1(\Omega \cap B(x_0, R))} < \exists \varepsilon_0 \Rightarrow x_0 \notin \mathcal{S}$$

nice cut-off function

$$\begin{array}{ll} \text{nice cut-off function} & x_0 \in \overline{\Omega}, \ 0 < R \ll 1, \\ \varphi = \varphi_{x_0,R} \in \mathcal{Y} \\ 1, \ x \in B(x_0, \frac{R}{2}) & |\nabla \varphi| \le CR^{-1}\varphi^{\frac{5}{6}} \\ 0, \ x \in \mathbf{R}^2 \setminus B(x_0, R) & |\nabla^2 \varphi| \le CR^{-2}\varphi^{\frac{2}{3}} & \mathcal{Y} = \{\varphi \in C^2(\overline{\Omega}) \mid \frac{\partial \varphi}{\partial \nu}\Big|_{\partial \Omega} = 0 \} \end{array}$$

Formation of collapses

symmetry of the Green function

$$\varphi \in C^2(\overline{\Omega}), \left. \frac{\partial \varphi}{\partial \nu} \right|_{\partial \Omega} = 0$$

weak form (symmetrization)

monotonicity formula  $\lambda = \|u(\cdot,t)\|_1$ 

$$\left|\frac{d}{dt}\int_{\Omega} u\varphi\right| \le C(\lambda + \lambda^2) \|\nabla\varphi\|_{C^1}$$

weak continuation

$$0 \leq \exists \mu(dx, t) \in C_*([0, T], \mathcal{M}(\overline{\Omega}))$$
$$u(x, t)dx = \mu(dx, t), \ 0 \leq t < T$$

arepsilon -regularity

$$\lim_{R \downarrow 0} \limsup_{t \uparrow T} \|u(\cdot, t)\|_{L^1(\Omega \cap B(x_0, R))} < \exists \varepsilon_0 \Rightarrow x_0 \notin \mathcal{S} \implies x_0 \in \mathcal{S} \Rightarrow \lim_{R \downarrow 0} \limsup_{t \uparrow T} \|u(\cdot, t)\|_{L^1(\Omega \cap B(x_0, R))} \ge \varepsilon_0$$

$$\longleftrightarrow \text{ monotonicity formula}$$

 $\lim_{R \downarrow 0} \liminf_{t \uparrow T} \|u(\cdot, t)\|_{L^1(\Omega \cap B(x_0, R)} \ge \varepsilon_0$ 

formation of collapse

$$\mu(dx,T) = \sum_{x_0 \in \mathcal{S}} m(x_0)\delta_{x_0} + f(x)dx, \ m(x_0) \ge \varepsilon_0, \ 0 \le f = f(x) \in L^1(\Omega)$$

Generation of the weak solution  

$$0 \leq \mu = \mu(dx, t) \in C_*([0, T], \mathcal{M}(\overline{\Omega})) \quad \text{weak solution}$$

$$\stackrel{\text{def}}{\longrightarrow} 0 \leq \exists \mathcal{N} = \mathcal{N}(\cdot, t) \in L^{\infty}_*([0, T], \mathcal{X}') \quad \text{multiplicated operator}$$
1.  $t \in [0, T] \mapsto \langle \varphi, \mu(dx, t) \rangle, \ \varphi \in \mathcal{Y} \quad \text{a.c.}$ 
2.  $\frac{d}{dt} \langle \varphi, \mu \rangle = \langle \Delta \varphi, \mu \rangle + \frac{1}{2} \langle \rho_{\varphi}, \mathcal{N}(\cdot, t) \rangle \text{ a.e. } t \in [0, T]$ 
3.  $\mathcal{N}|_{C(\overline{\Omega} \times \overline{\Omega})} = \mu \otimes \mu$ 

$$\begin{array}{ll} \textbf{Theorem 1} & \mu_k(dx,t) \in C_*([0,T],\mathcal{M}(\overline{\Omega})) \\ & \mathcal{N}_k \in L^\infty_*([0,T],\mathcal{X}') \ \text{ weak solutions} \end{array}$$

 $\begin{array}{l} 0 \leq \mu_k(\overline{\Omega}, t) \leq C \\ \|\mathcal{N}_k(\cdot, t)\|_{\mathcal{X}'} \leq C \end{array} \quad \longrightarrow \quad \text{sub-sequence} \end{array}$ 

$$\begin{split} \mu_k(dx,t) &\rightharpoonup \mu(dx,t) \quad \text{in } C_*([0,T],\mathcal{M}(\overline{\Omega})) \\ \mathcal{N}_k(\cdot,t) &\rightharpoonup \mathcal{N}(\cdot,t) \quad \text{in } L^\infty_*([0,T],\mathcal{X}') \quad \text{weal solution} \end{split}$$

$$\mathcal{Y} = \{ \varphi \in C^2(\overline{\Omega}) \mid \frac{\partial \varphi}{\partial \nu} \Big|_{\partial \Omega} = 0 \} \quad \mathcal{X} = [\mathcal{X}_0]^{L^{\infty}(\Omega \times \Omega)}$$
$$\mathcal{X}_0 = \{ \rho_{\varphi} + \psi \mid \varphi \in \mathcal{Y}, \ \psi \in C(\overline{\Omega} \times \overline{\Omega}) \}$$

$$\stackrel{\bullet}{\longrightarrow} \quad \mu(\overline{\Omega}, t) = \mu(\overline{\Omega}, 0) \equiv \lambda, \ 0 \le t \le T \\ \left| \frac{d}{dt} \langle \varphi, \mu(dx, t) \rangle \right| \le C(\lambda + \lambda^2) \| \nabla \varphi \|_{C^1}$$

 $u = u(x,t) \quad \text{classical solution}$   $\implies \mathcal{N}(\cdot,t) = u(x,t) \otimes u(x',t) \ dxdx'$   $\|\mathcal{N}(\cdot,t)\|_{\mathcal{X}'} = \lambda^2, \ \lambda = \|u_0\|_1$ 

Backward self-similar transformation
$$x_0 \in S$$
 $y = (x - x_0)/(T - t)^{1/2}, \ s = -\log(T - t)$  $z(y, s) = (T - t)u(x, t)$ weak limit $s_k \uparrow +\infty$  subsequence $z(y, s + s_k)dy \rightarrow \exists \zeta(dy, s)$  in  $C_*(-\infty, +\infty; \mathcal{M}(\mathbf{R}))$ Put $z(y, s) = 0$  where it is not defined

 $(2^{2}))$ Put  $\mathcal{A}(g, s)$ U where it is not defined.

 $\mathcal{M}(\mathbf{R}^2) = C_\infty(\mathbf{R}^2)'$  $C_{\infty}(\mathbf{R}^2) = \{ z \in C(\mathbf{R}^2 \cup \{\infty\}), \ z(\infty) = 0 \}$ 

$$\begin{aligned} \left| \frac{d}{dt} \int_{\Omega} u(\cdot, t) \varphi_{x_0, R} \right| &\leq C_{\lambda} R^{-2}, \ 0 < R \ll 1 \\ \left| \langle \varphi_{x_0, R}, u(\cdot, t) dx \rangle - \langle \varphi_{x_0, R}, \mu(dx, T) \rangle \right| &\leq C_{\lambda} (T - t) / R^2 \\ s_k + s &= -\log(T - t), \ R = b(T - t)^{1/2} \\ &\longrightarrow \\ \left| \langle \varphi_{0, b}, z(\cdot, s + s_k) dy \rangle - \langle \varphi_{x_0, be^{-(s + s_k)/2}}, \mu(dx, T) \rangle \right| \\ &\leq C_{\lambda} / b^2 \\ \mu(dx, T) &= \sum_{x_0 \in \mathcal{S}} m(x_0) \delta_{x_0}(dx) + f(x) dx \\ k \to \infty, \ b \uparrow + \infty \quad \longrightarrow \quad \boxed{m(x_0) = \zeta(\mathbf{R}^2, s)} \end{aligned}$$

$$u(x,t)dx \rightharpoonup \sum_{x_0 \in \mathcal{S}} m(x_0)\delta_{x_0}(dx) + f(x)dx$$

Second parabolic envelope

$$\left\langle |y|^2, \zeta(dy,s) \right\rangle \leq C$$

 $\zeta_s = \nabla \cdot (\nabla \zeta - \zeta \nabla (\Gamma * \zeta + |y|^2/4))$  in  $\mathbf{R}^2 \times (-\infty, +\infty)$ 

$$\begin{aligned} & \operatorname{Proof} \qquad \varphi \in C_0^2(\mathbf{R}^2), \, s \gg 1 \\ & \frac{d}{ds} \int_{\mathcal{O}_s} z\varphi = \int_{\mathcal{O}_s} (\partial_s \varphi + y \cdot \nabla \varphi + \Delta \varphi) z \\ & + \frac{1}{2} \int_{\mathcal{O}_s \times \mathcal{O}_s} \rho_{\varphi}^s(y, y') z \otimes z \\ & \mathcal{O}_s = \Omega_s \times \{s\} \\ & \rho_{\varphi}^s(y, y') = \nabla \varphi(y) \cdot \nabla_y G_s(y, y') \\ & + \nabla \varphi(y') \cdot \nabla_{y'} G_s(y, y') \end{aligned}$$
$$\begin{aligned} & G(x, x') = \Gamma(x - x') + K(x, x') \\ & (x, x') \in (\overline{\Omega} \times \Omega) \cup (\Omega \times \overline{\Omega}) \\ & \Gamma(x) = \frac{1}{2\pi} \log \frac{1}{|x|} \\ & G_s(y, y') = \Gamma(y - y') - \frac{s}{4\pi} \\ & + K(e^{-s}y + x_0, e^{-s}y' + x_0) \end{aligned}$$

Exclusion of boundary blowup 
$$x_0 \in \partial\Omega$$
  
 $0 \leq \zeta(dy, s), \ \zeta(\mathbf{R}^2, s) \leq \lambda \equiv ||u_0||_1, \ \text{supp } \zeta(dy, s) \subset \overline{\mathbf{R}^2_+}$   
 $\zeta_s = \nabla \cdot (\nabla\zeta - \zeta\nabla(E * \zeta + |y|^2/4)) \ \text{in } \mathbf{R}^2_+ \times (-\infty, +\infty)$   
 $\frac{\partial z}{\partial \nu} - z \frac{\partial}{\partial \nu} (E * \zeta + \frac{|y|^2}{4}) \Big|_{\partial \mathbf{R}^2_+} = 0$ 

 $E(y, y') = \Gamma(y - y') - \Gamma(y - y'_*)$ 



$$\begin{split} \varphi &= |y|^2 \psi_R, \, \psi_R(y) = \psi(y/R) \\ \psi &= \varphi_{0,2}(|y|) \\ \Delta \varphi &= 4\psi_R + 4\frac{y}{R} \cdot \nabla \psi(\frac{y}{R}) + \frac{|y|^2}{R^2} \Delta \psi(\frac{y}{R}) \\ y \cdot \nabla \varphi &= 2|y|^2 \psi_R + |y|^2 \frac{y}{R} \cdot \nabla \psi(\frac{y}{R}) \\ \langle 1 + |y|^2, \zeta(dy, s) \rangle \leq C \end{split}$$

 $\Rightarrow (\text{dominated convergence theorem})$  $\lim_{R\uparrow+\infty} \int_{s_1}^{s_2} \langle \Delta \varphi + \frac{y}{2} \cdot \nabla \varphi, \zeta(dy, s) \rangle$  $= 4(s_2 - s_1)m(x_0) + \int_{s_1}^{s_2} I(s)ds$  $I(s) = \langle |y|^2, \zeta(dy, s) \rangle$ 

 $\frac{dI}{ds} = 4m(x_0) + I(s) \text{ a.e. } s \Rightarrow \lim_{R \uparrow +\infty} I(s) = +\infty$ 

Proof of Theorem B (continued)  $x_0 \in \mathcal{S}$ 

#### backward self-similar transformation

$$y = (x - x_0)/(T - t)^{1/2}, \ s = -\log(T - t)$$
$$z(y, s) = (T - t)u(x, t)$$

weak limit 
$$s_k \uparrow +\infty$$
 subsequence

$$z(y, s+s_k)dy \rightharpoonup \exists \zeta(dy, s) \text{ in } C_*(-\infty, +\infty; \mathcal{M}(\mathbf{R}^2))$$

limit equation exclusion of boundary blowup 
$$x_0 \in \Omega$$
  
 $\zeta_s = \nabla \cdot (\nabla \zeta - \zeta \nabla (\Gamma * \zeta + |y|^2/4))$  in  $\mathbf{R}^2 \times (-\infty, +\infty)$ 

$$u(x,t)dx \rightharpoonup \sum_{x_0 \in \mathcal{S}} m(x_0)\delta_{x_0}(dx) + f(x)dx$$

## parabolic envelope

$$m(x_0) = \zeta(\mathbf{R}^2, s) \qquad \left\langle |y|^2, \zeta(dy, s) \right\rangle \le C$$

## scaling back

$$\zeta(dy,s) = e^{-s} A(dy',s'), \ y' = e^{-s/2}y, \ s' = -e^{-s}$$

$$A_s = \nabla \cdot (\nabla A - A \nabla \Gamma * A) \text{ in } \mathbf{R}^2 \times (-\infty, 0)$$
$$A = A(dy, s) \ge 0, \ A(\mathbf{R}^2, s) = m(x_0)$$

$$\begin{aligned} \epsilon - \text{regularity} & \rightarrow \zeta^s(dy, s) = \sum_{j=1}^{m(s)} \tilde{m}_j(s) \delta_{y_j(s)}(dy) \end{aligned} \qquad A^s(w_j(s)) \leq m(s) \leq m(x_0) / \varepsilon_0, \ |y_j(s)| \leq C, \ \tilde{m}_j(s) \geq \varepsilon_0 \end{aligned}$$

singular part 
$$A^s(dy',s') = \sum_{j=1}^{m(s')} \tilde{m}_j(s') \delta_{y'_j(s)}(dy')$$

$$A_s = \nabla \cdot (\nabla A - A \nabla \Gamma * A) \text{ in } \mathbf{R}^2 \times (-\infty, 0)$$
$$A = A(dy, s) \ge 0, \ A(\mathbf{R}^2, s) = m(x_0)$$

singular part
$$A^s(dy',s') = \sum_{j=1}^{m(s')} ilde{m}_j(s') \delta_{y_j'(s)}(dy')$$

scaling limit 
$$s'_0 < 0, \ 1 \le j \le m(s'_0)$$
 fix  
 $\tilde{A}_{\beta}(dy', s') = \beta^2 A(dy, s)$   
 $y = \beta y' + y'_j(s'_0), \ s = \beta^2 s' + s'_0$ 

 $\tilde{A}_{\beta_k}(dy', s') \rightharpoonup \tilde{A}(dy', s') \in C_*(-\infty, s'_0; \mathcal{M}(\mathbf{R}^2))$ 

 $eta_k \downarrow 0$  subsequence

 $ilde{A}(dy',s')=m_{i}'(s_{0}')\delta_{0}(dy')$  weak solution

#### translation limit

$$\begin{split} s'_k \uparrow +\infty, \ \hat{A}_k(dy', s') &= \tilde{A}(dy', s' + s'_k)_{\text{subsequence}} \\ \hat{A}_k(dy', s') &\rightharpoonup \hat{A}(dy', s') \in C_*(-\infty, +\infty; \mathcal{M}(\mathbf{R}^2)) \\ \hat{A}(dy', s') &= m'_j(s'_0)\delta_0(dy') \end{split}$$

weak Liouville propertyweak solution (measure valued) $a_s = \nabla \cdot (\nabla a - a \nabla \Gamma * a)$  in  $\mathbf{R}^2 \times (-\infty, +\infty)$  $\Rightarrow a(\mathbf{R}^2, s) = 0$  or  $8\pi$ Kurokiba-Ogawa03

$$\longrightarrow \tilde{m}_j(s'_0) = \tilde{A}(\mathbf{R}^2, 0) = 8\pi$$

if 
$$A^{ac}(dy',s') = 0 \longrightarrow A(dy',s') = 8\pi \sum_{j=1}^{\ell} \delta_{y'_j(s')}(dy')$$

 $ightarrow m(x_0)=8\pi\ell$  collapse mass quantization

# $\exists (\varepsilon_0, R) > 0$

Gagliardo-Nirenberg

$$\longrightarrow \frac{dJ}{dt} + 3J^{3/2} \le C_R, \ J = \int_{\Omega} u(\log u - 1) + 1 \ dx$$
$$\frac{d}{dt} t^{-2} + 3(t^{-2})^{3/2} = t^{-3}$$
$$J(t) \le t^{-2}, \ 0 < t \le \min\{t_1, t_0\}, \ t_0^{-3} = C_R \quad \square$$
parabolic regularity

scaling

Theorem 3 
$$\exists \varepsilon_0, \sigma_0, C$$
  
 $u_t = \Delta u - \nabla \cdot (u \nabla \Gamma * u) \text{ in } \mathbf{R}^2 \times (-T, T)$   
 $u_0 = u|_{t=0}$  weak solution generated by classical solutions  
 $\|u_0\|_{L^1(B(x_0,2R))} < \varepsilon_0, \ u_0 = u|_{t=0} \Rightarrow$   
 $\sup_{t \in [-\sigma_0 R^2, \sigma_0 R^2] \cap (-T,T)} \|u(\cdot,t)\|_{L^{\infty}(B(x_0,R))} \leq CR^{-2}$ 

scaling invariant regularity (inverse scaling back)  

$$\zeta(B(y_0, 2r), s) < \varepsilon_0 \implies \|\zeta(\cdot, s)\|_{L^{\infty}(B(y_0, r))} \le Cr^{-2}$$

Lemma 
$$u_t = \Delta u - \nabla \cdot (u \nabla \Gamma * u)$$
 in  $\mathbf{R}^2 \times (0, T)$   
 $u|_{t=0} = u_0(x) \ge 0$ 

$$\exists \varepsilon_0, \ R, \ t_0 > 0, \ \forall x_0 \qquad \|u_0\|_{L^1(B(x_0, 8R))} < \varepsilon_0/2$$

$$\forall \tau \in (0, t_0)$$
  
$$\sup_{\tau \le t < t_0} \|u(\cdot, t)\|_{L^{\infty}(B(x_0, R))} < +\infty$$

**Proof** 
$$\exists t_1 \in (0,T), \|u_0\|_{L^1(B(x_0,8R))} < \varepsilon_0/2$$

$$\longrightarrow \sup_{t \in (0,t_1)} \|u(\cdot,t)\|_{L^1(B(x_0,4R))} < \varepsilon_0$$

$$\rightarrow \frac{d}{dt} \int_{\mathbf{R}^2} u(\log u - 1)\varphi \ dx + \frac{1}{8} \int_{\mathbf{R}^2} u^{-1} |\nabla u|^2 \varphi \ dx \\ \leq C_{\varphi}, \ 0 \leq t < t_1, \ \varphi = \varphi_{x_0,R}$$

 $\|u(\cdot,t)\|_{L\log L(B(x_0,2R))}$  Involved by the initial value!

residual vanishing

ng 
$$A^{ac}(dy',s') = 0 \iff \zeta^{ac}(dy,s) = 0$$

1<sup>st</sup> envelope $m(x_0) = \zeta(\mathbf{R}^2, s)$ 

$$\begin{array}{l} \mathbf{2^{nd}\ envelope} \\ \left\langle |y|^2, \zeta(dy,s) \right\rangle \leq C \end{array}$$

 $\zeta_{s} = \nabla \cdot \left( \nabla \zeta - \zeta \nabla (\Gamma * \zeta + |y|^{2}/4) \right)$ scaling invariant regularity attractive potential toward infinity  $\zeta(dy, s)$   $s \in \mathbf{R}$ 



outer second moment

$$\frac{d}{ds}\langle\varphi,\zeta\rangle \geq \langle\Delta\varphi - C\varphi_r + \frac{1}{2}r\varphi_r,\zeta\rangle, \ \varphi = \varphi(r)$$

$$\varphi(r) = \xi(r/R), \ \xi(r) = r^{-} - 1$$
$$R \gg 1 \implies \Delta \varphi + \frac{1}{2}r\varphi_r \ge C\varphi_r, \ r \ge R$$

$$\frac{d}{ds}\langle (\frac{|y|^2}{R^2} - 1)_+, \zeta(dy, s) \rangle \ge 0 \quad \longrightarrow \quad \zeta(dy, s) = \zeta^s(dy, s)$$

uniform estimate of self-interaction part

improved regularity  $\zeta(B(y_0, 2r), s) < \varepsilon_0$   $\longrightarrow \|\zeta(\cdot, s)\|_{L^{\infty}(B(y_0, r))} \le Cr^{-2}$ 



Tracing by the local second moment

#### simple blowup point

$$\ell = 1 \implies \zeta(dy, s) = 8\pi\delta_0(dy)$$





Senba-S. 01	weak formulation monotonicity formula	formation of collapse weak solution generation instant blowup for over mass concentrated initial data non-existence of over mass entire solution without concentration
Senba-S. 02a	weak solution	
Kurokiba-Ogawa 03	scaling invariance	
Senba-S. 04	backward self-similar transformation scaling limit	
S. 05	parabolic envelope (1) scaling invariance of the scaling limit a local second moment	sub-collapse quantization
		collapse mass quantization
Senba-Ohtsuka-S. 07	defect moment (1)	radially symmetric dynamics
Senba 07, Naito-S. 08	parabolic envelope (2)	type II blowup rate
S. 08	scaling back	limit equation simplification
Senba-S. 11 (2 <sup>nd</sup> ed.)	translation limit	concentration-cancelation simplification
S. 13a	limit equation classification	boundary blowup exclusion
S. 13b	improved regularity concentration compactness	cloud formation collision of sub-collapses
S. 14	tightness	residual vanishing
S. 18	defect moment (2)	quantization of BUIT
S. 22	outer second moment	residual vanishing in finite time